

Les fluctuations de courant : au delà du bruit

B. Reulet (1,2), D. Prober (1), J. Gabelli (2)

(1) Yale University – Department of Applied Physics (USA)

(2) Laboratoire de Physique des Solides – Orsay (France)

Introduction:

what to measure ?

Averaged quantities

1) STATIC (thermodynamic or steady state) properties:

$\vec{M}(\vec{H})$ Magnetization vs. Magnetic field

$I(V)$ Current vs. Voltage

$\vec{n}(\vec{E})$ Molecular orientation vs. Electric field

$x(F)$ Displacement vs. Force

X ray diffraction vs. V,P, etc...

Susceptibilities

2) DYNAMICAL properties: response to an oscillating field

$$H = H_{dc} + \delta H \cos \omega_0 t \quad \chi(H_{dc}, \omega_0) = \frac{\partial M_{\omega_0}}{\partial H_{\omega_0}} \quad \text{Magnetic susceptibility}$$
$$V = V_{dc} + \delta V \cos \omega_0 t \quad G(V_{dc}, \omega_0) = \frac{\partial I_{\omega_0}}{\partial V_{\omega_0}} \quad \text{ac Conductance}$$

Both are COMPLEX quantities: in-phase and out-of-phase response.

Interesting when:

$$\omega_0 \tau > 1$$

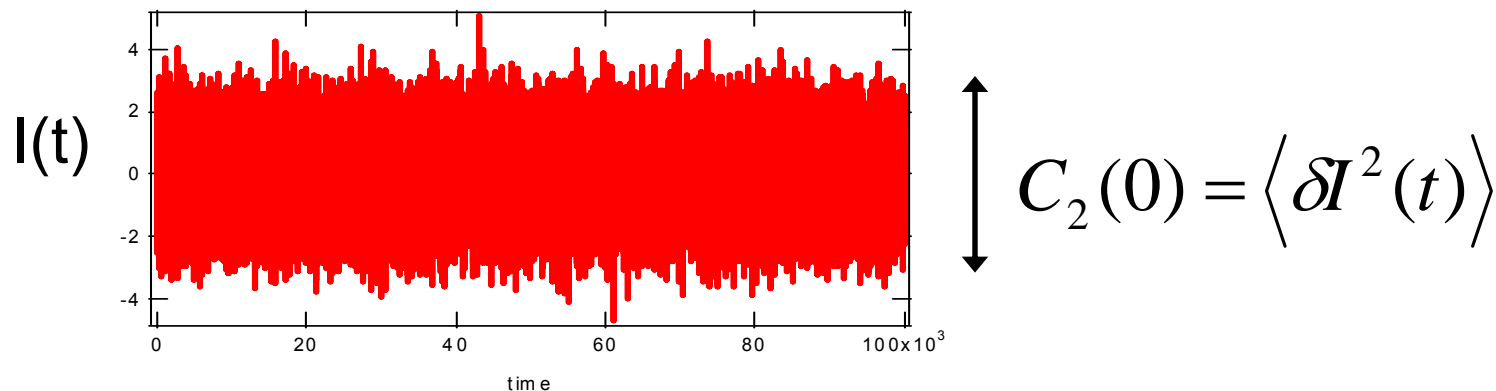
Aging of polymers or spin glasses, dielectric constant or magnetic permittivity, etc.

Fluctuations (noise)

3) Intrinsic FLUCTUATIONS (or noise):

Correlation function: $C_2(V, \tau) = \langle \delta I(t + \tau) \delta I(t) \rangle$ Time average ↙

Noise spectral density: $S_2(V, \omega) = \langle \delta I(\omega) \delta I(-\omega) \rangle$



Magnetization noise (motion of domain walls), $1/f$ noise of electric dipoles in amorphous systems glasses, of voltage in transistors, electric noise associated with depinning of charge density waves, etc.

Summary

Average

$$\langle \bullet \rangle$$

Fluctuations

$$\langle \bullet \bullet \rangle$$

$$I(V)$$

$$S_2(V, \omega) = \langle \delta I(\omega) \delta I(-\omega) \rangle$$

Dynamical
response

$$\frac{\partial \bullet}{\partial V_{\omega_0}}$$

$$G(V, \omega_0) = \frac{\partial I_{\omega_0}}{\partial V_{\omega_0}}$$

Beyond

	$\langle \bullet \rangle$	$\langle \bullet \bullet \rangle$	$\langle \bullet \bullet \bullet \rangle$
	$I(V)$	$\langle \delta I(\omega) \delta I(-\omega) \rangle$	$\langle \delta I(\omega) \delta I(\omega') \delta I(-\omega - \omega') \rangle$
$\frac{\partial \bullet}{\partial V_{\omega_0}}$	$\frac{\partial I_{\omega_0}}{\partial V_{\omega_0}}$	$\frac{\partial S_2(\omega)}{\partial V_{\omega_0}}$	<p>THIRD CUMULANT</p> <p>↑</p>
$\frac{\partial^2 \bullet}{\partial V_{\omega_0} \partial V_{\omega_1}}$	$\frac{\partial^2 I_{\omega_0 \pm \omega_1}}{\partial V_{\omega_0} \partial V_{\omega_1}}$		

NOISE DYNAMICS

Frequency mixing (non linear transport)

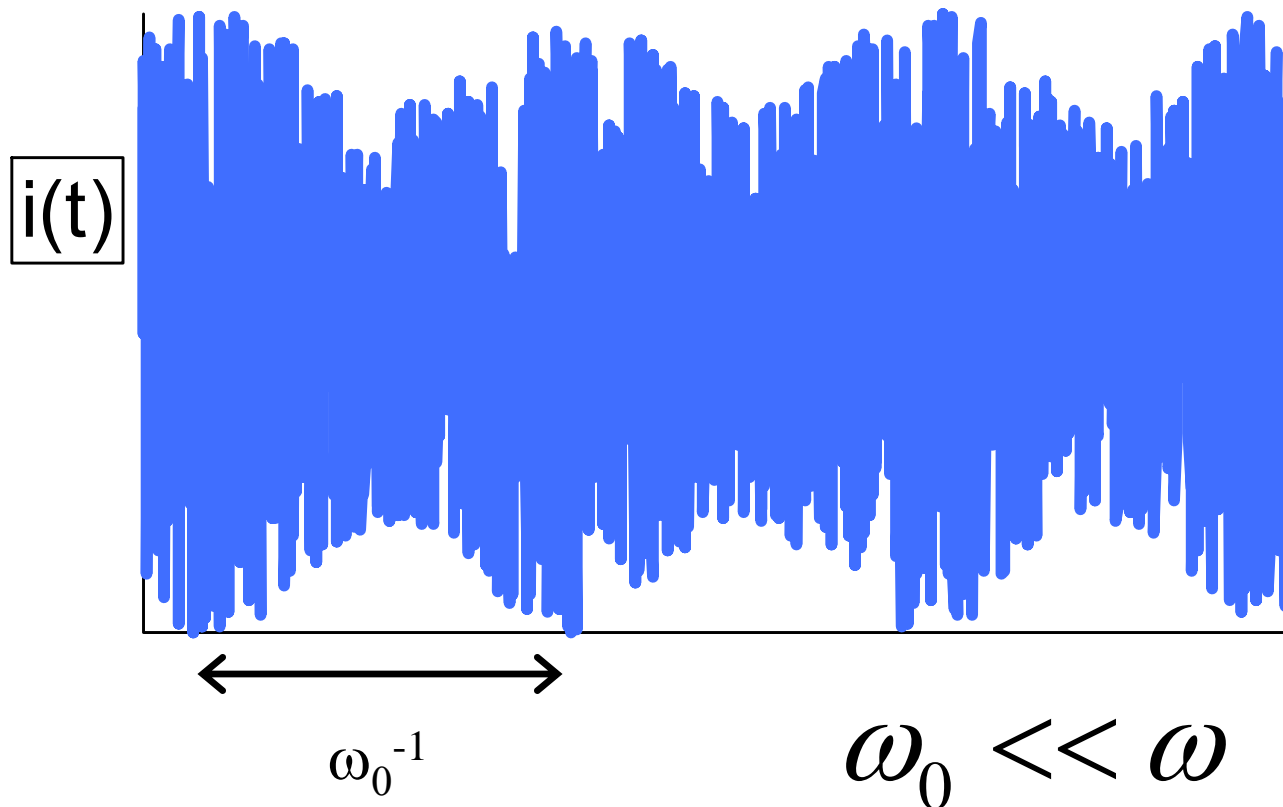
The Dynamics of noise:

The noise susceptibility
in the
classical regime

Noise susceptibility – How fast can one modulate noise ?

$$V(t) = V_{dc} + \delta V \cos \omega_0 t$$

$$\frac{\partial S_2(\omega)}{\partial V_{\omega_0}}$$



Noise susceptibility – the case of a macroscopic conductor

Fluctuation-dissipation theorem, at equilibrium and low frequency $\hbar\omega \ll k_B T$:

$$S_2 = 4k_B T G$$

NOISE = electron THERMOMETER

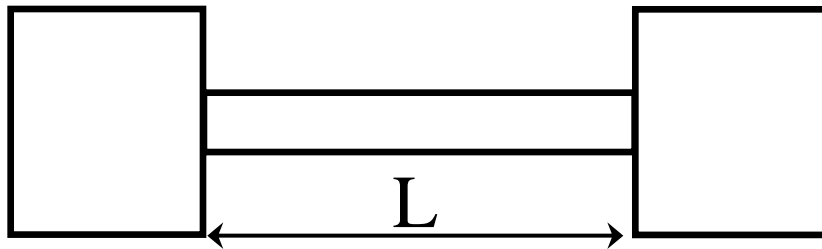
$$\delta V(t) \Rightarrow \delta P_{Joule}(t) = I \delta V(t) \Rightarrow \delta T(t) \Rightarrow \delta S_2(t)$$

$$\chi_{\omega_0}(\omega) = \frac{\partial S_2(\omega)}{\partial V_{\omega_0}} \propto Z_{thermal}^{-1}(\omega_0) \quad \omega \gg \omega_0$$

Its frequency dependence gives ENERGY RELAXATION (i.e. INELASTIC) time

Noise susceptibility – from macro- to mesoscopic conductor

diffusive metallic wire: length $L \gg$ mean free path



$$L^2 = D \tau_D$$

* long wire or SNS: phonon cooling

The noise susceptibility gives the ELECTRON-PHONON time

* intermediate wire: diffusion cooling

The noise susceptibility gives the DIFFUSION time

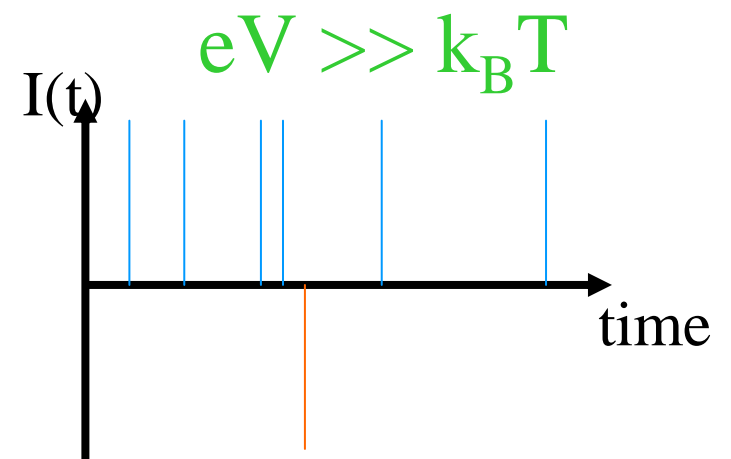
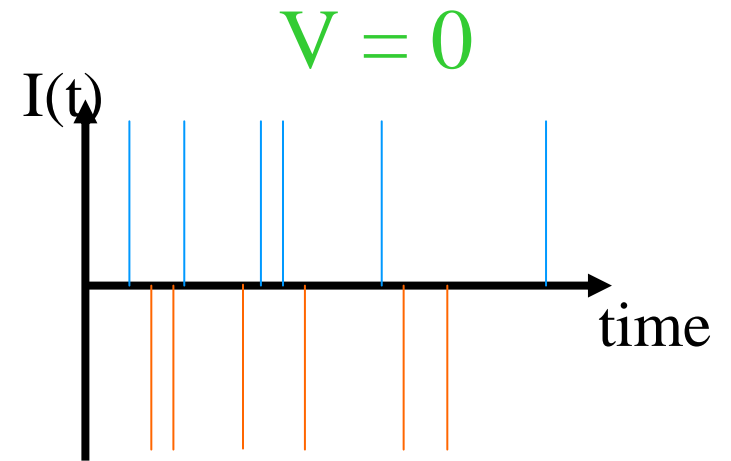
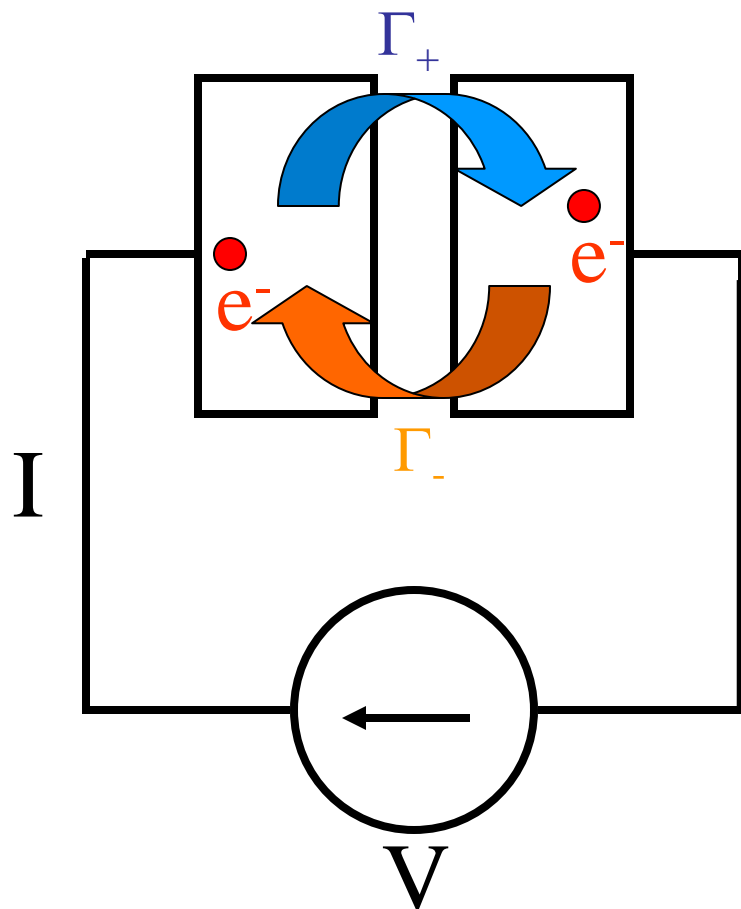
* short wire: elastic transport (independent electrons)

The noise susceptibility gives the ELECTRON-ELECTRON time

* ballistic wire (nanotube), quasi-crystals (sub-diffusive), ... ??

The noise susceptibility
in the
Quantum regime

Current in a tunnel junction



Current fluctuations in a tunnel junction at low frequency

$$\langle \delta I^2 \rangle = 2eIB \coth\left(\frac{eV}{2k_B T}\right) \quad \text{B=bandwidth}$$

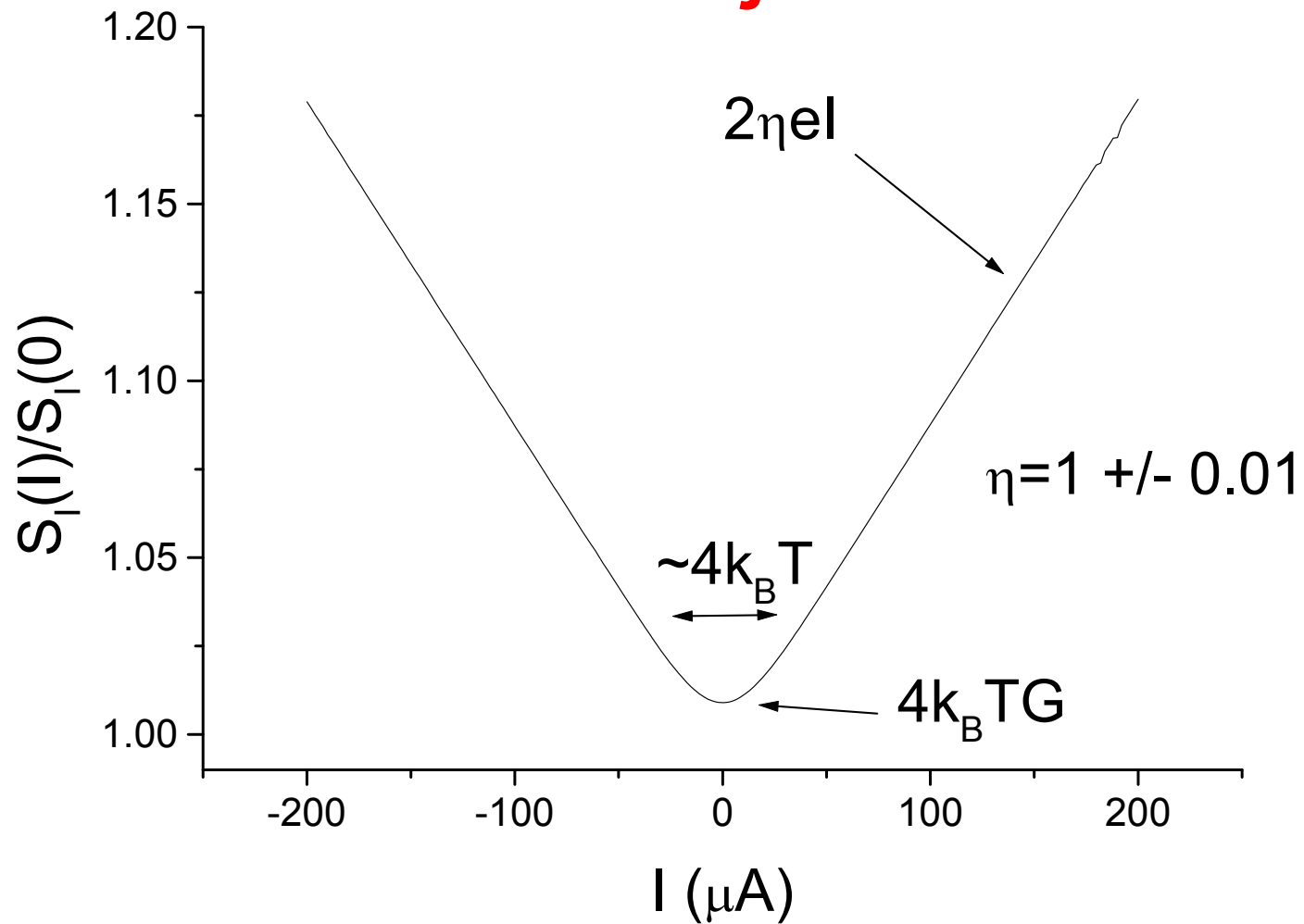
$$S_2 = \begin{cases} 4k_B T G & \text{if } eV \ll k_B T \\ 2eI & \text{if } eV \gg k_B T \end{cases}$$

↑
Noise spectral density in A²/Hz

← Equilibrium (Johnson) noise:
macroscopic, fluctuation-
dissipation theorem

← Shot noise: discreteness of
charge

Measurement of S_2 on a tunnel junction



Quantum mechanics: ordering of operators?

Average current:

$$I_{DC} = \langle \hat{I} \rangle$$

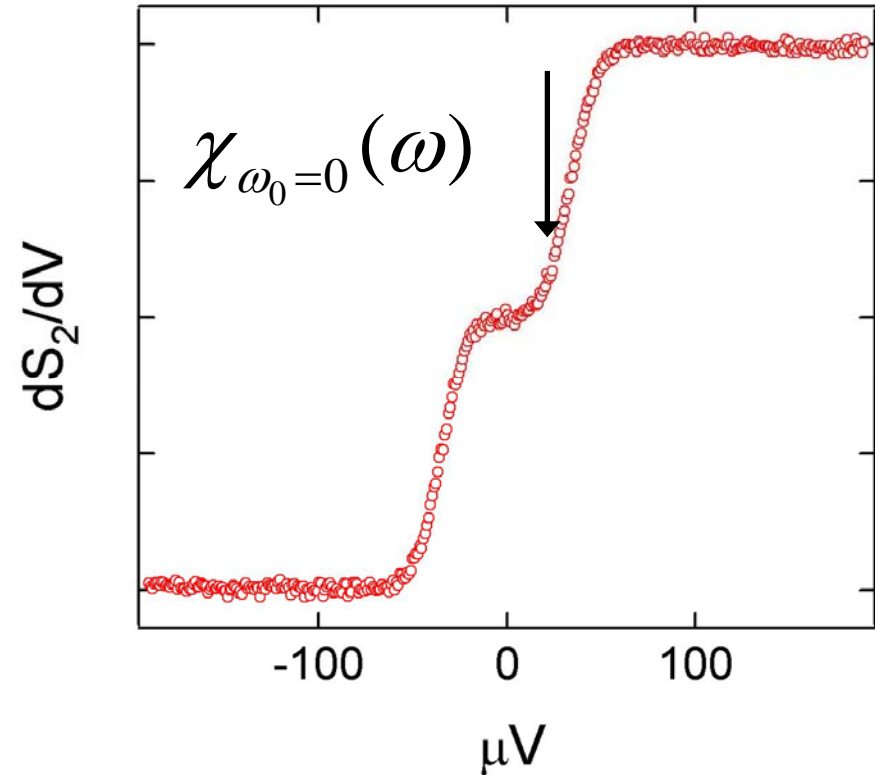
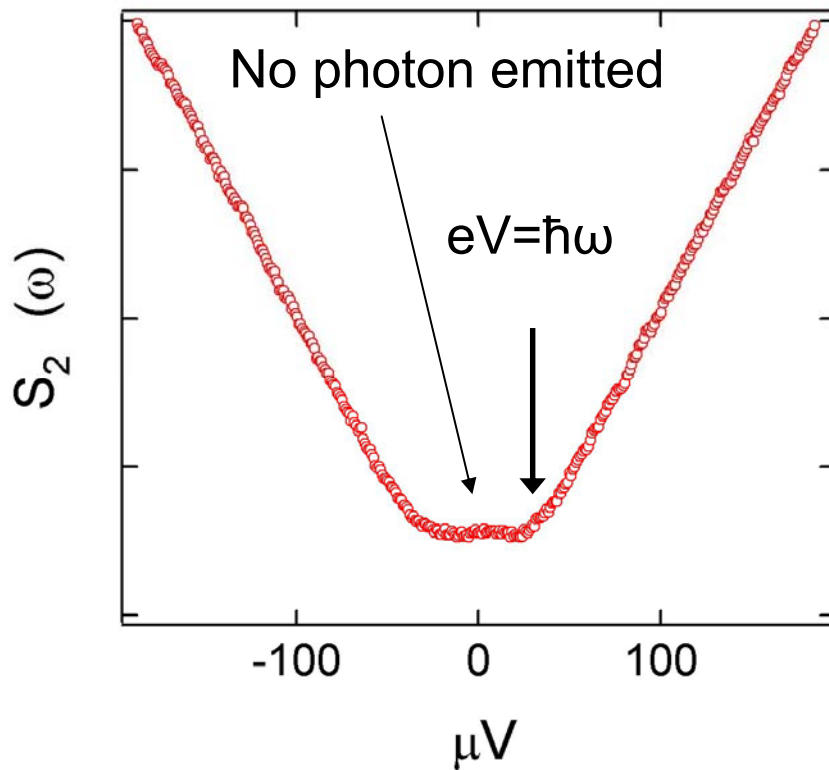
Noise S_2 :

$$S_2(\omega) = \int dt e^{i\omega t} \left\{ \begin{array}{l} \langle \hat{I}(0) \hat{I}(t) \rangle \\ \langle \hat{I}(t) \hat{I}(0) \rangle \\ \frac{1}{2} \left(\langle \hat{I}(0) \hat{I}(t) \rangle + \langle \hat{I}(t) \hat{I}(0) \rangle \right) \end{array} \right. \begin{array}{l} \text{Absorption} \\ \text{Emission} \\ \text{Classical} \end{array}$$

$$S_2^{abs}(\omega) = S_2^{em}(-\omega)$$

$$S_2^{sym}(\omega) = S_2^{em}(\omega) + \frac{1}{2} G \hbar \omega$$

Reminder: noise in the quantum regime $\hbar\omega > k_B T, eV$ for a tunnel junction



Measured on a tunnel junction at $T=35\text{mK}$, $f=6\text{ GHz}$, $\hbar\omega/k_B T \sim 8.5$

Noise susceptibility – beyond the classical regime: theory

What if $\omega_0 > \omega$?
What if $\hbar\omega > k_B T$?

$$\chi_{\omega_0}(\omega) \propto \langle \delta I(\omega) \delta I(\omega_0 - \omega) \rangle$$

Calculation:

- * Landauer-Büttiker formalism
- * SYMMETRIZATION of the operators and $\omega \rightarrow -\omega$

The symmetrization rule depends on the experimental setup !

In particular: $\omega \sim \omega_0$

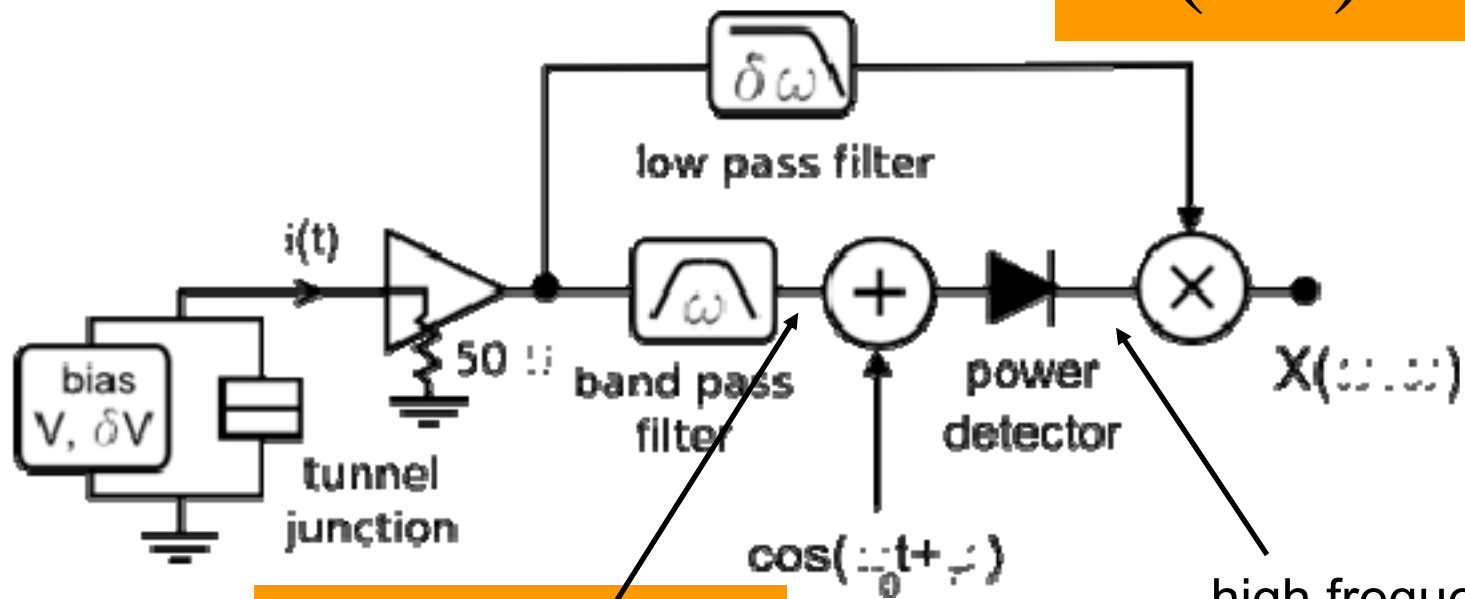
$$\chi_{\omega}(\omega) \propto \langle \delta I(\omega) \delta I(0) \rangle$$

Noise susceptibility – the quantum regime: experiment

$\omega_0 \sim \omega \sim 6 \text{ GHz}$
 $\hbar\omega/k_B T \sim 8.5$
 $\delta\omega \sim 100 \text{ MHz}$

low frequency current

$$\delta I(\pm \varepsilon) e^{\pm i \varepsilon t}$$

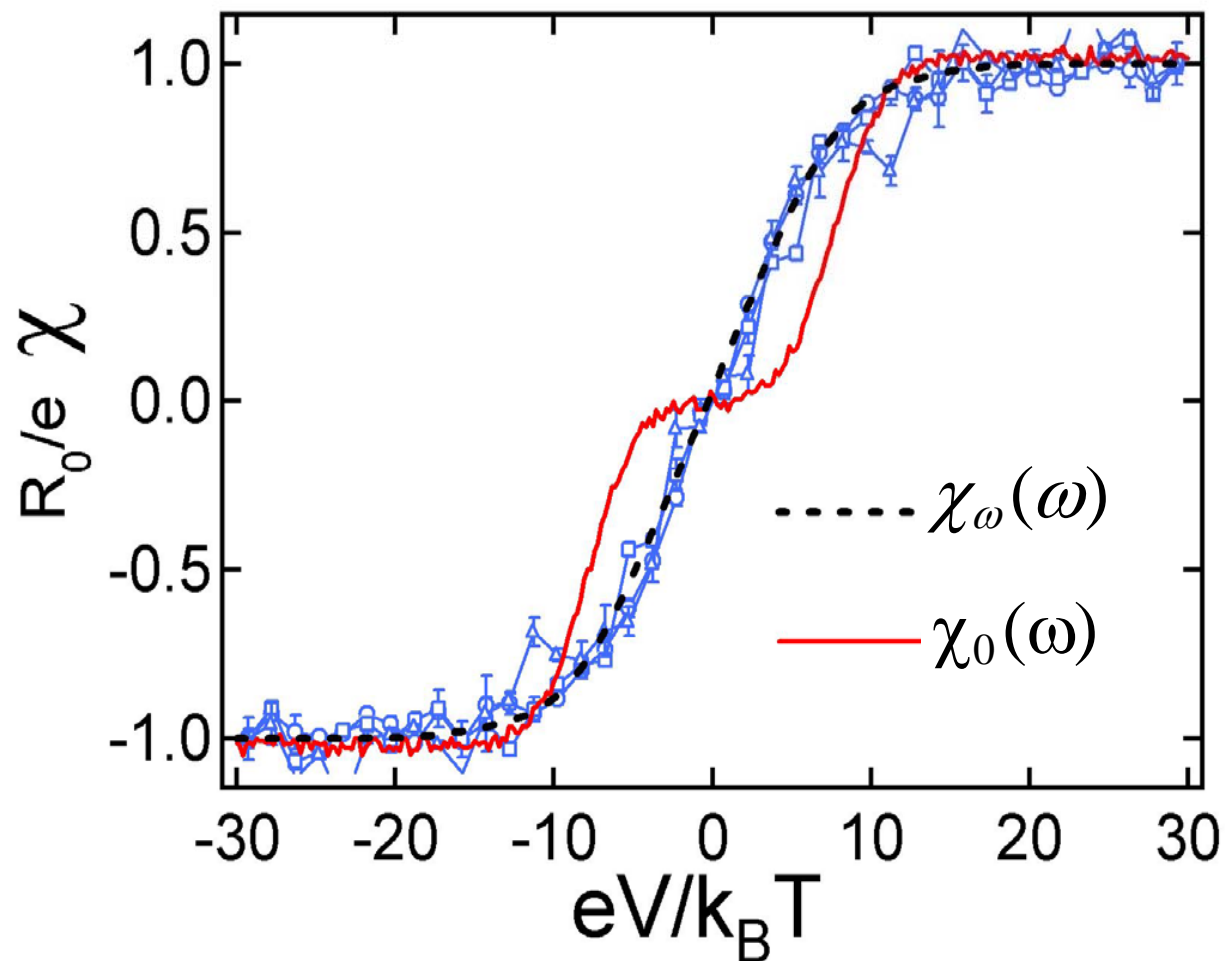


$$\delta I(\pm \omega) e^{\pm i \omega t}$$

high frequency current shifted to low freq.

Noise susceptibility – the quantum regime: experiment

$\omega_0 \sim \omega \sim 6$ GHz
 $T = 35$ mK
 $\hbar\omega/k_B T \sim 8.5$

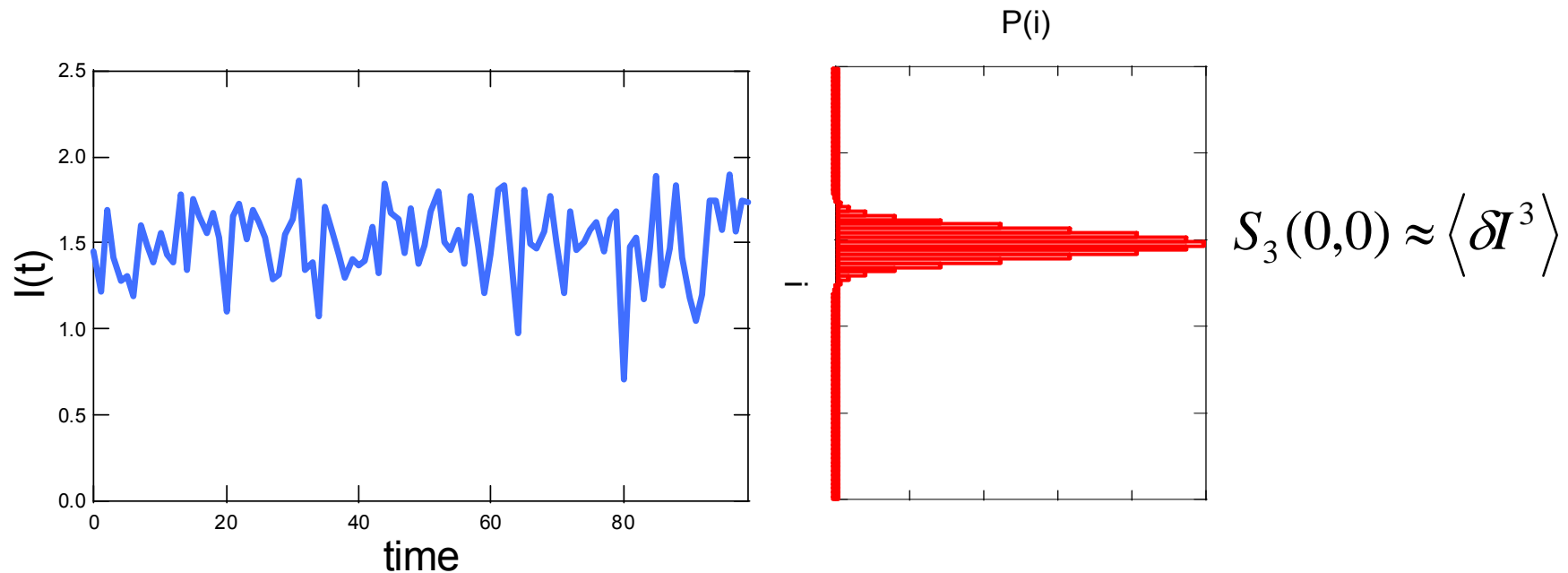


The third cumulant
of current fluctuations
in the classical regime

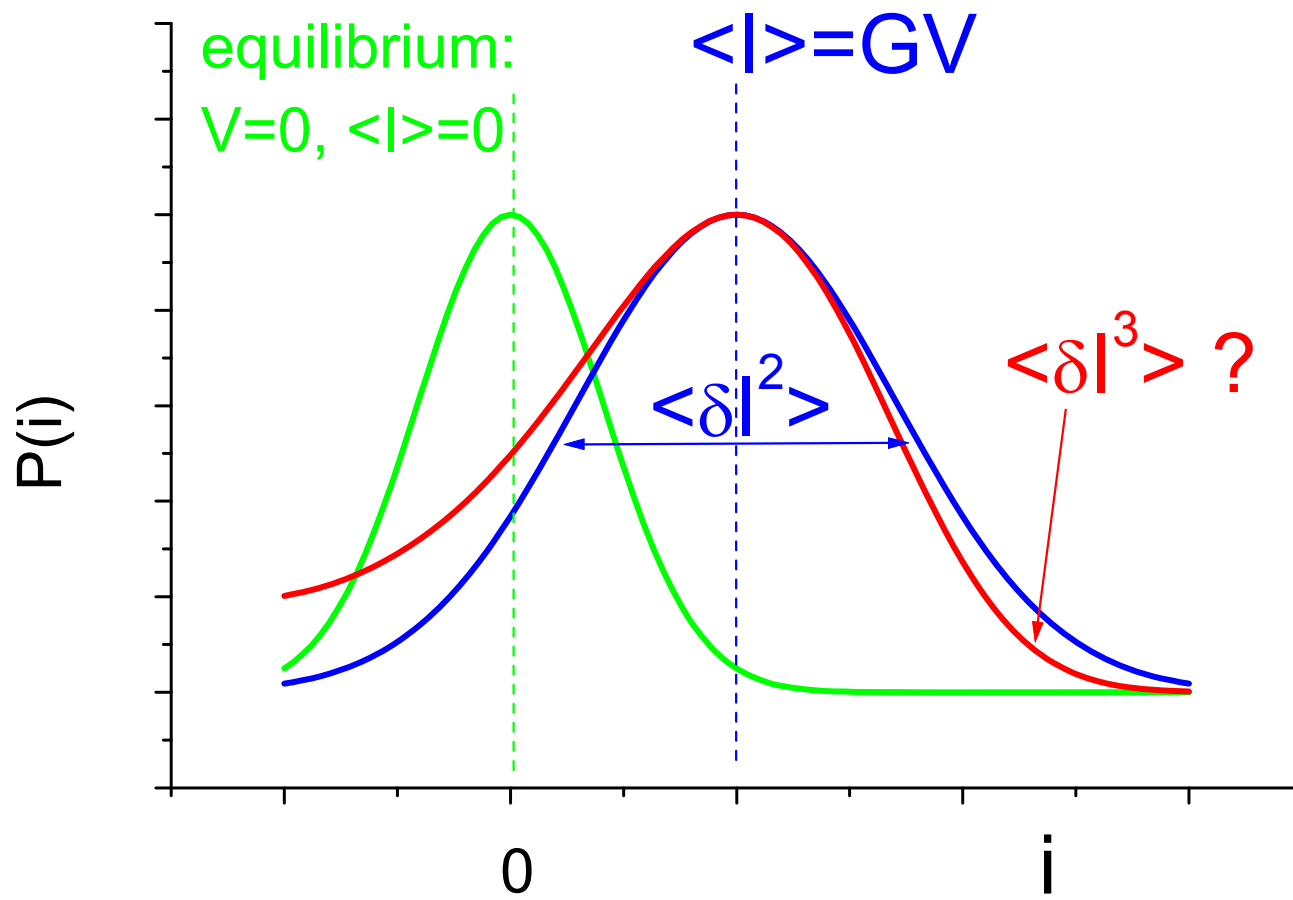
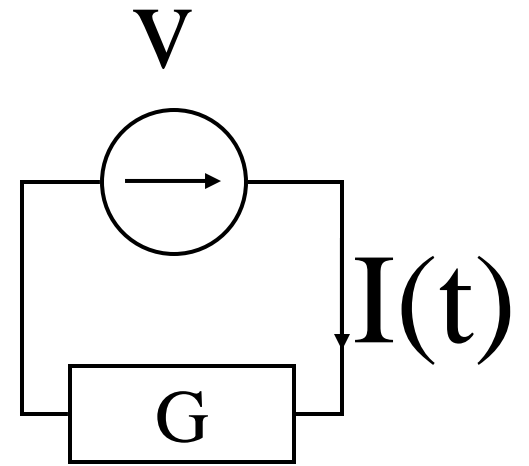
The third cumulant of noise

$$S_3(\omega, \omega') = \langle \delta I(\omega) \delta I(\omega') \delta I(-\omega - \omega') \rangle$$

At low frequency: $S_3 =$ **SKEWNESS** of the probability distribution of current fluctuations $P(I)$: **zero for gaussian noise**



Statistics of the current

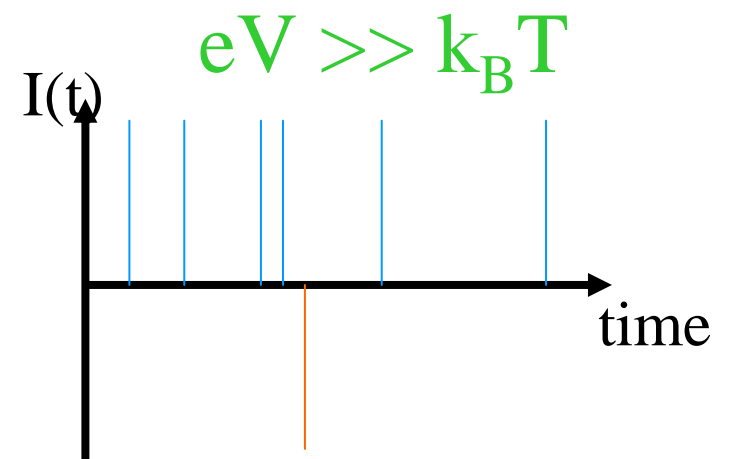
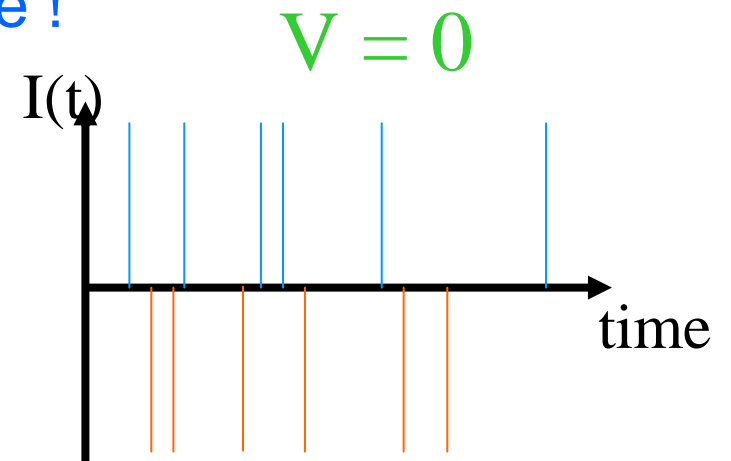
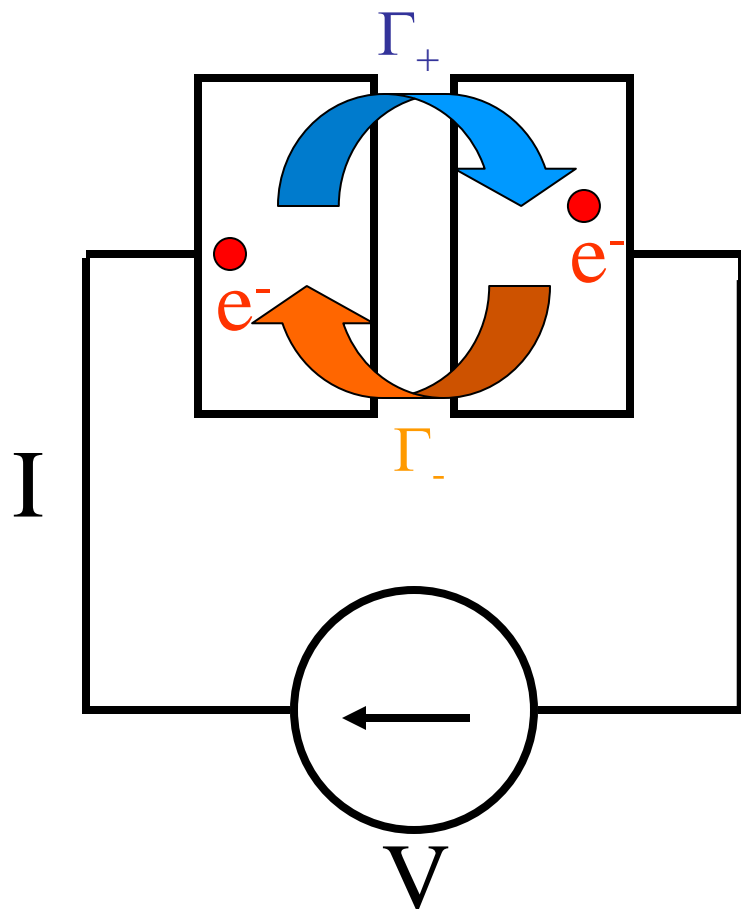


$$I(t) = \langle I \rangle + \delta I(t)$$

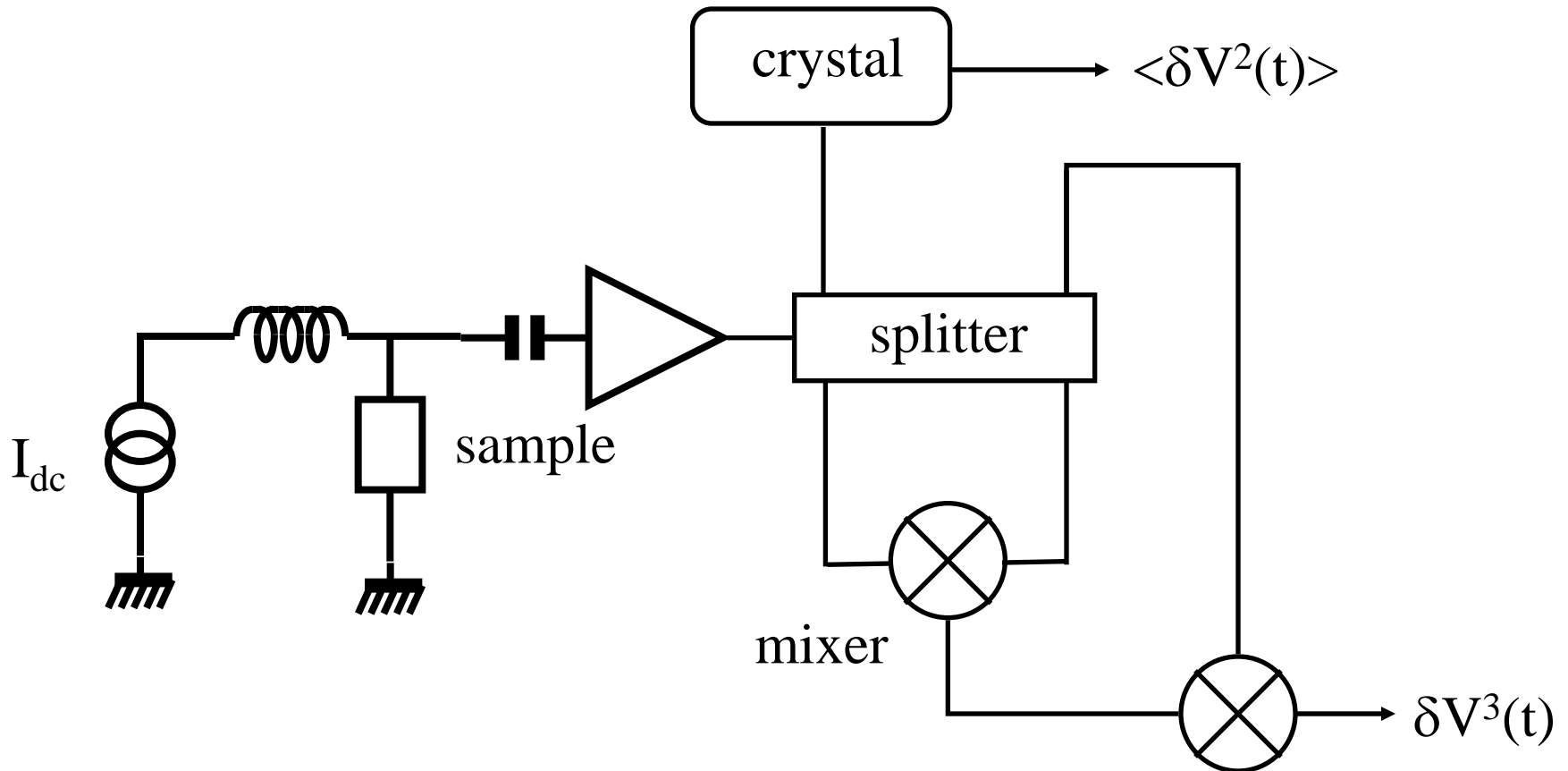
At equilibrium: Johnson noise: $\langle \delta I^2 \rangle = 4k_B TGB$, (B=bandwidth)

Current in a tunnel junction

$S_3 = e^2 I$ independent of temperature !

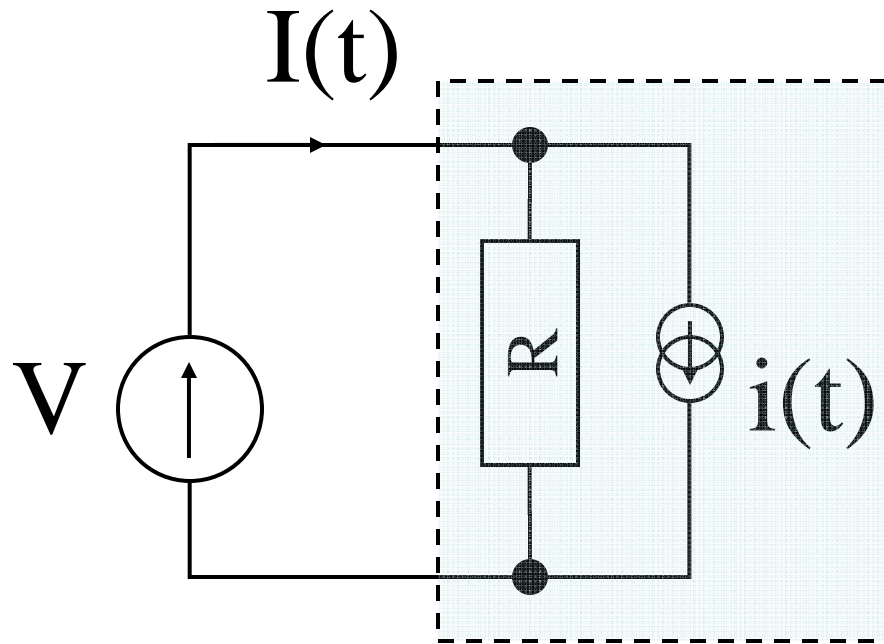


Measurement: method

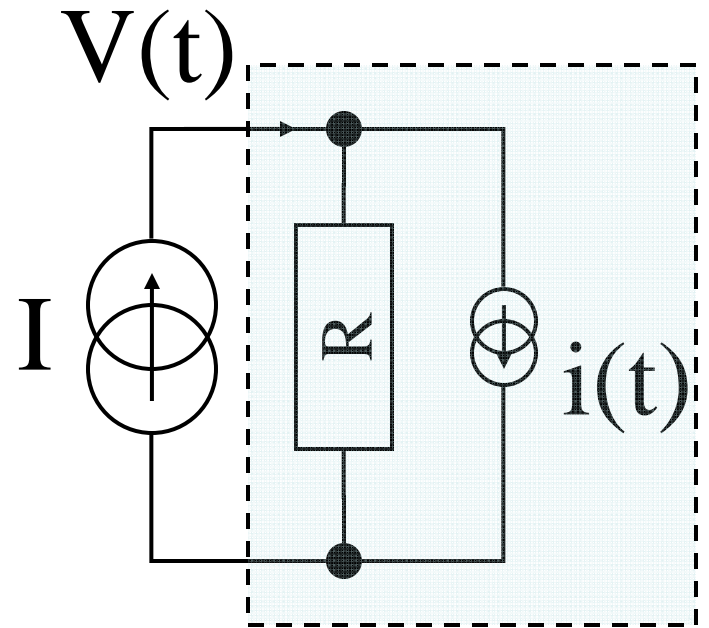


Bandwidth: 10-1200 MHz

Voltage bias vs. Current bias



$$\delta I(t) = i(t)$$



$$\delta V(t) = -Ri(t)$$

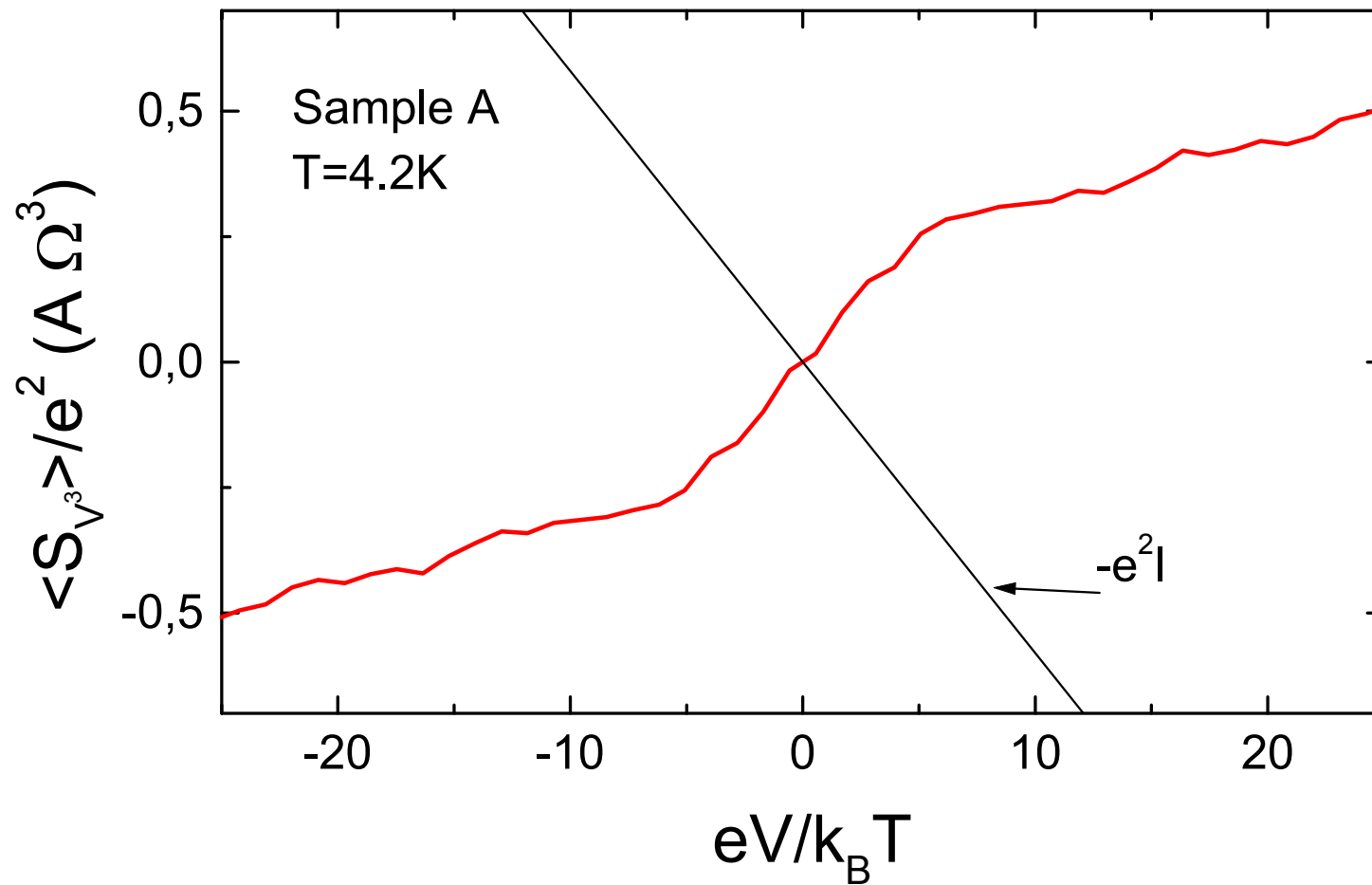
$$V = RI$$

$$\langle \delta V^2 \rangle = R^2 \langle \delta I^2 \rangle$$

$$\langle \delta V^3 \rangle = -R^3 \langle \delta I^3 \rangle ?$$

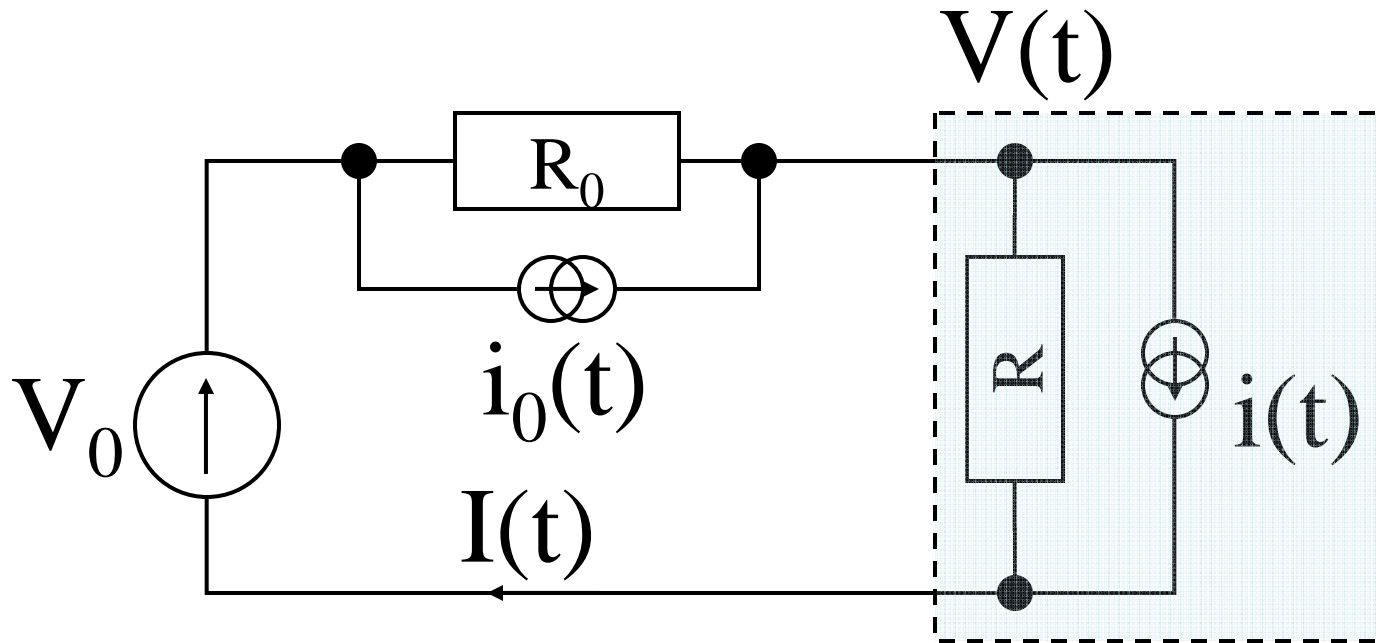
Result (T=4.2K)

(averaged over 12 days !)



Environmental effects

Imperfect voltage bias



$$\delta V(t) = (R // R_0) (i_0(t) - i(t))$$

External noise Feedback of the environment

The probability distribution $P(i)$ depends on $V(t)$

Effect of the environment on the probability distribution $P(i)$

$$P(i, V) = P_V(i, V - R_0 i) \cong P_V(i, V) - R_0 i \frac{\partial P_V(i, V)}{\partial V}$$

Probability when voltage biased

Effect on the moments:

$$\langle i^n \rangle = \langle i^n \rangle_V - R_0 \frac{\partial \langle i^{n+1} \rangle_V}{\partial V}$$

Origin: e-e interactions

Application: n=3

* Effect on the third moment:

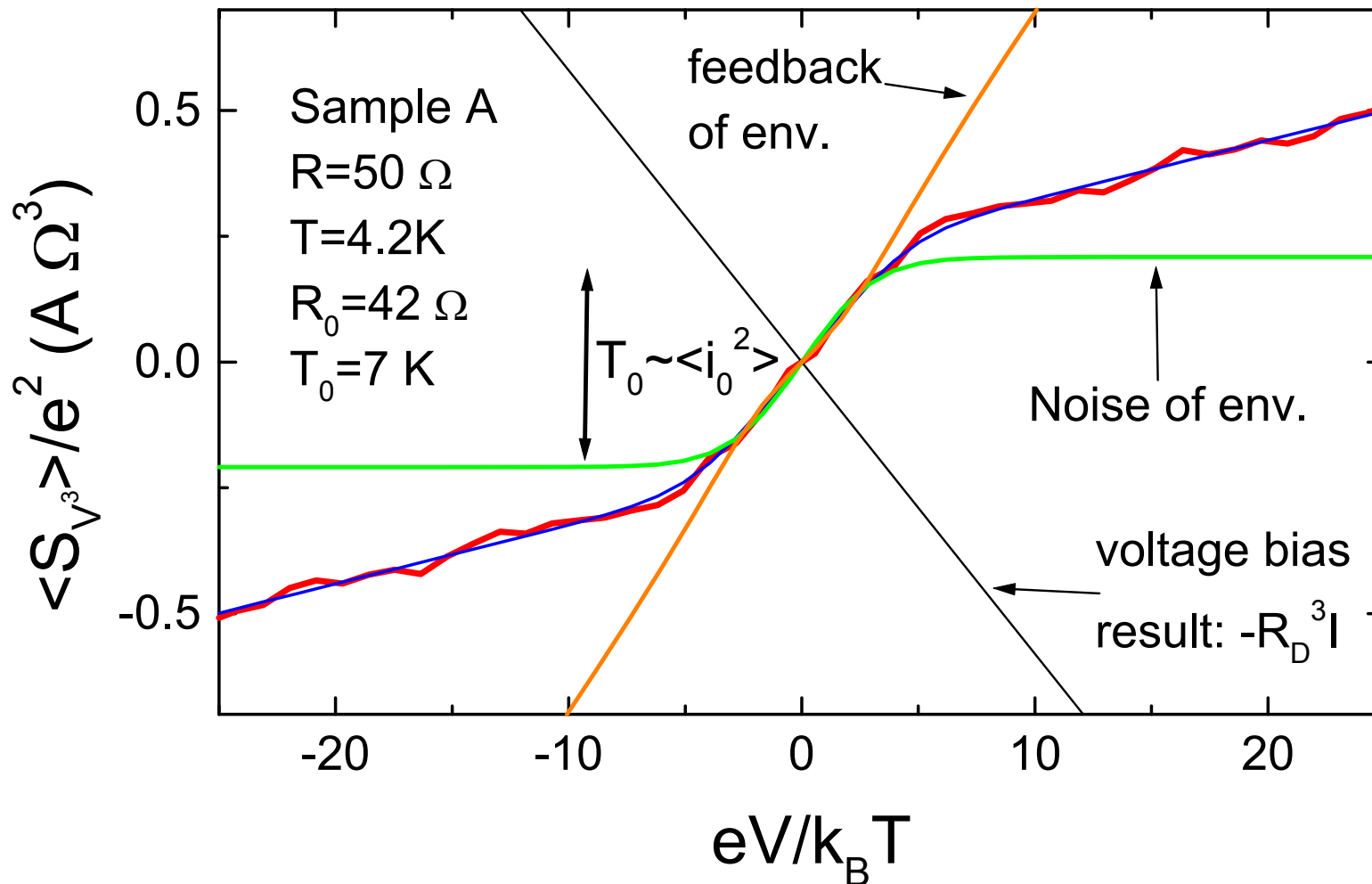
$$\langle i^3 \rangle = e^2 I - R_0 \frac{\partial S_4}{\partial V} \approx e^2 I - 6R_0 S_2 \frac{\partial S_2}{\partial V}$$



Noise Susceptibility
at $\omega_0 \sim 0$

$$\chi_{\omega_0=0}(\omega) = \frac{\partial S_2(\omega)}{\partial V}$$

Exp. Result + theory



Effect of the environment on the probability distribution $P(i)$

$$\langle i^n \rangle = \langle i^n \rangle_V - R_0 \frac{\partial \langle i^{n+1} \rangle_V}{\partial V}$$

Application: $n=1$

* Effect on the dc current:
dynamical Coulomb blockade

$$\langle i \rangle = I_{dc} - R_0 \frac{\partial S_2}{\partial V}$$

Noise Susceptibility
at $\omega_0 \sim \omega$

$$\chi_\omega(\omega)$$

Link between Dynamical Coulomb Blockade and Noise Susceptibility

$$\delta I = \frac{\hbar}{e^2} \int K(-\omega) \omega \chi_{\omega}(\omega) d\omega$$

Correlation function of
the environment

Noise susceptibility for $\omega_0 = \omega$

cf Lamb shift...

For small environmental impedance:

$$\delta I = - \int \text{Re} Z_0(\omega) \chi_{\omega}(\omega) d\omega$$

There is nothing quantum ...

The third cumulant
of current fluctuations
in the quantum regime

S_3 and Q mechanics: ordering ???

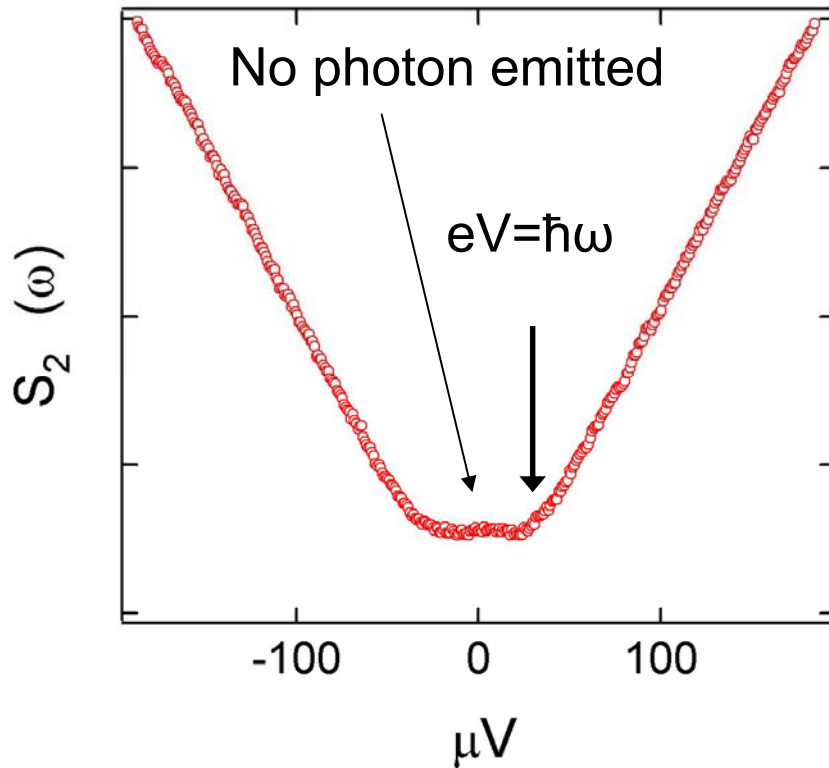
$$S_3(\omega, \omega') = \int dt dt' e^{i(\omega t + \omega' t')} \langle \hat{I}(0, t, t') \hat{I}(0, t, t') \hat{I}(0, t, t') \rangle$$

Two results:

$$S_3(0,0) = \frac{e^2}{h} V \cdot \begin{cases} t(1-t)(1-2t) & \text{Keldysh ordering} \\ t^2(1-t) & \text{Fully symmetrized} \end{cases}$$

$t \approx 10^{-5}$ for a tunnel junction

The third cumulant in the quantum regime: $S_3(0, \omega)$ with $\hbar\omega > k_B T, eV$

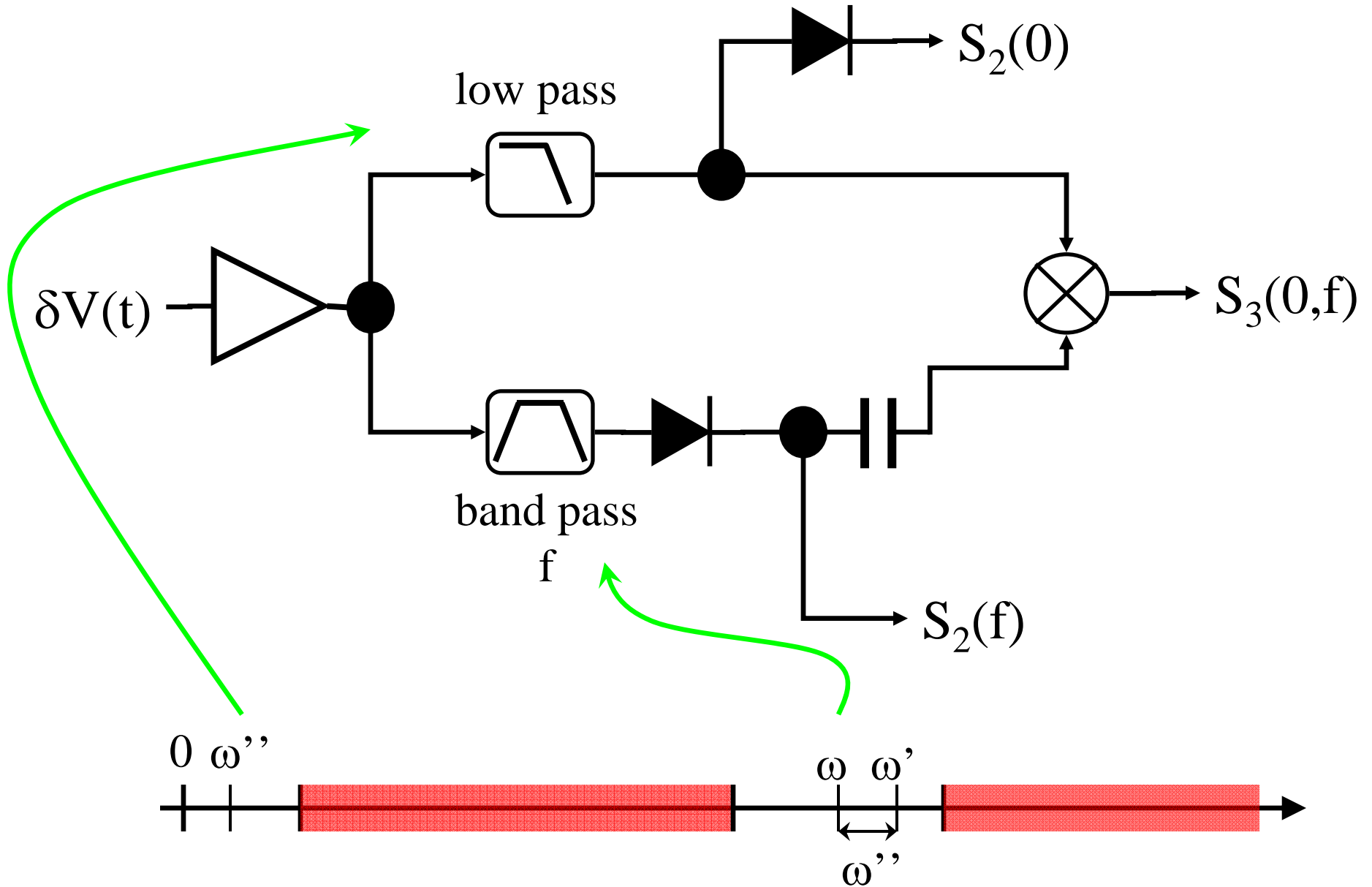


Prediction:

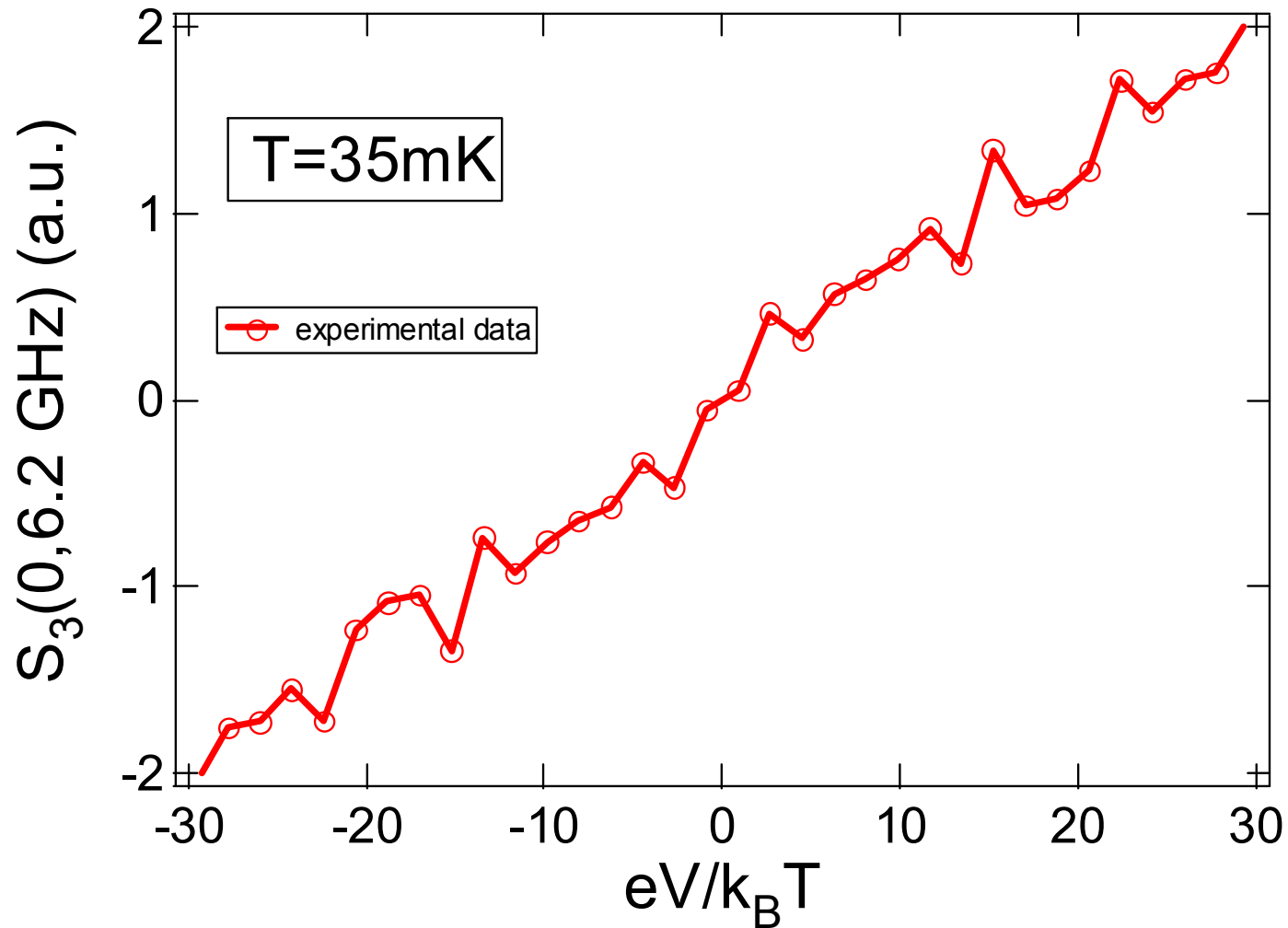
$$S_3(\omega, \omega') = e^2 I$$

indep of frequency !!

How to measure $S_3(0, f)$?



$S_3(0, \omega)$: experiment - minus environment in the Q regime



The third cumulant in the quantum regime: $S_3(0, \omega)$ with $\hbar\omega > k_B T, eV$

$$S_2(\omega) = \langle \delta I(\omega) \delta I(-\omega) \rangle \quad \sim \text{power of light emitted at frequency } \omega$$

=0 for $\hbar\omega > eV$

$$S_3(0, \omega) = \langle \delta I(0) \delta I(\omega) \delta I(-\omega) \rangle \quad \sim \text{correlations between low frequency current fluctuations and power of light at frequency } \omega.$$

Prediction:

$$S_3(\omega, \omega') = e^2 I \text{ independent of frequency !?}$$

WARNING: zero point motion of electrons !!

Another way to measure $S_3(0, f)$?

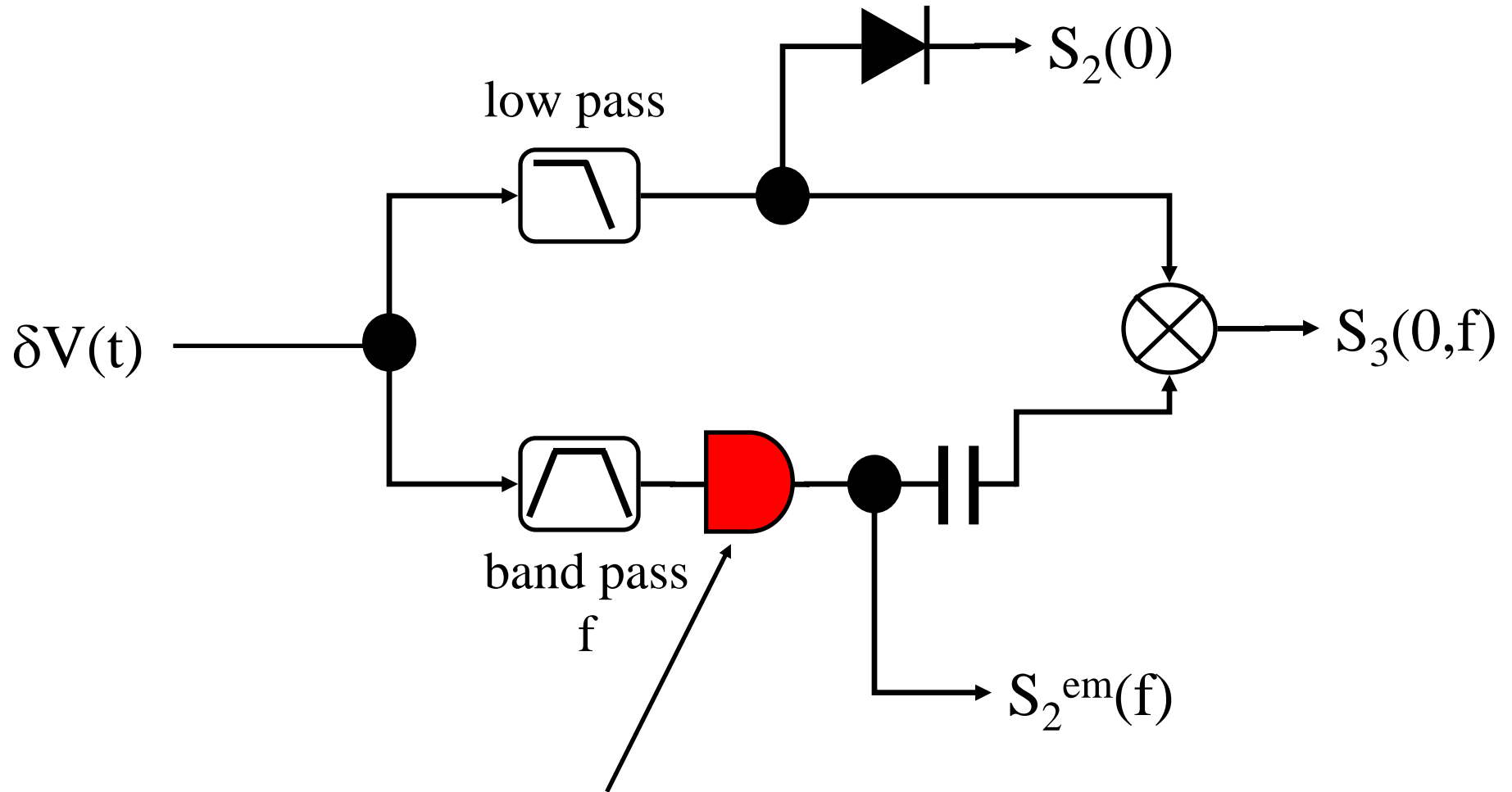
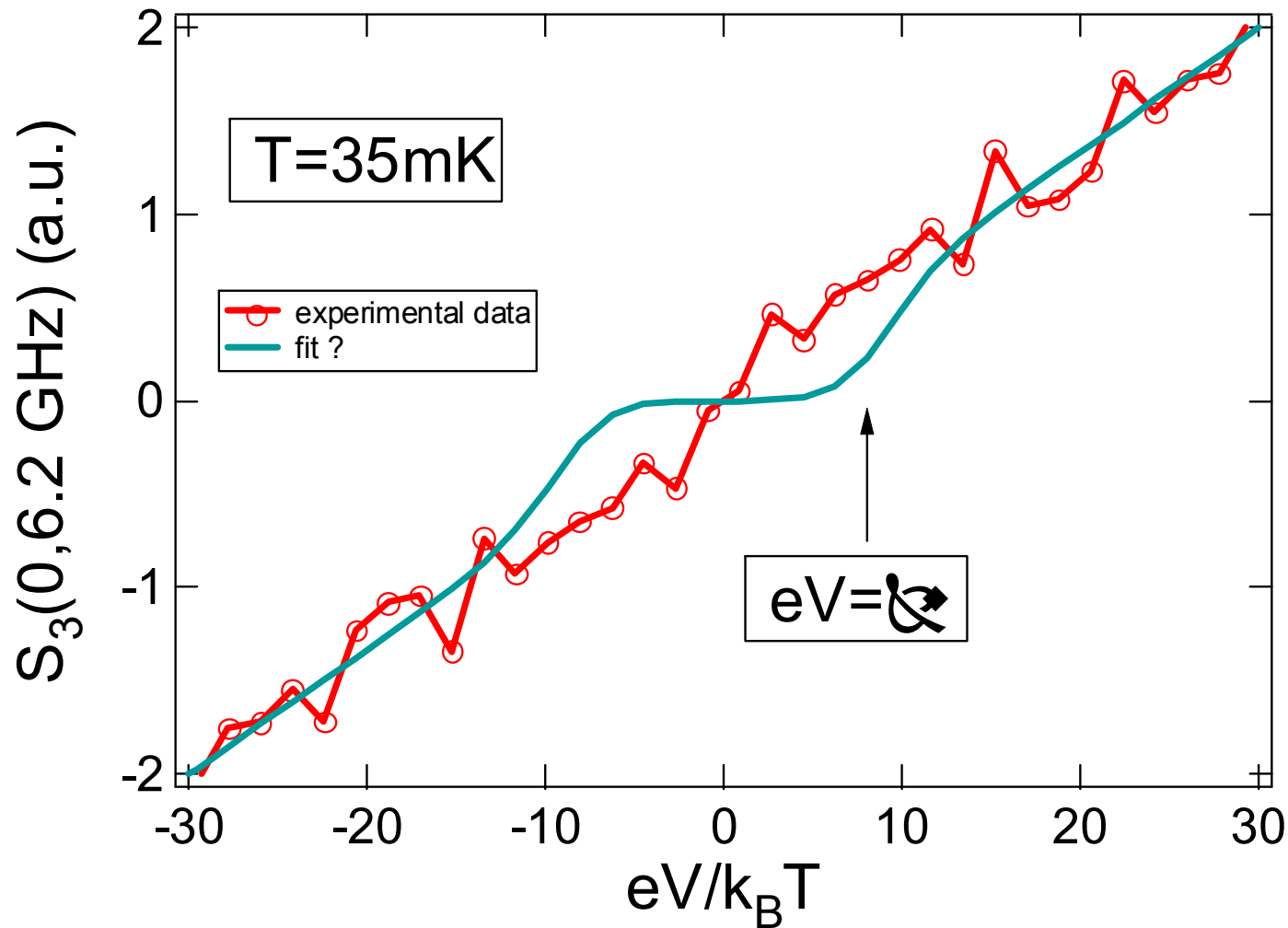


Photo-multiplier: absorbs photons

$S_3(0, \omega)$: experiment - minus environment in the Q regime



Another way to measure $S_3(0,f)$?

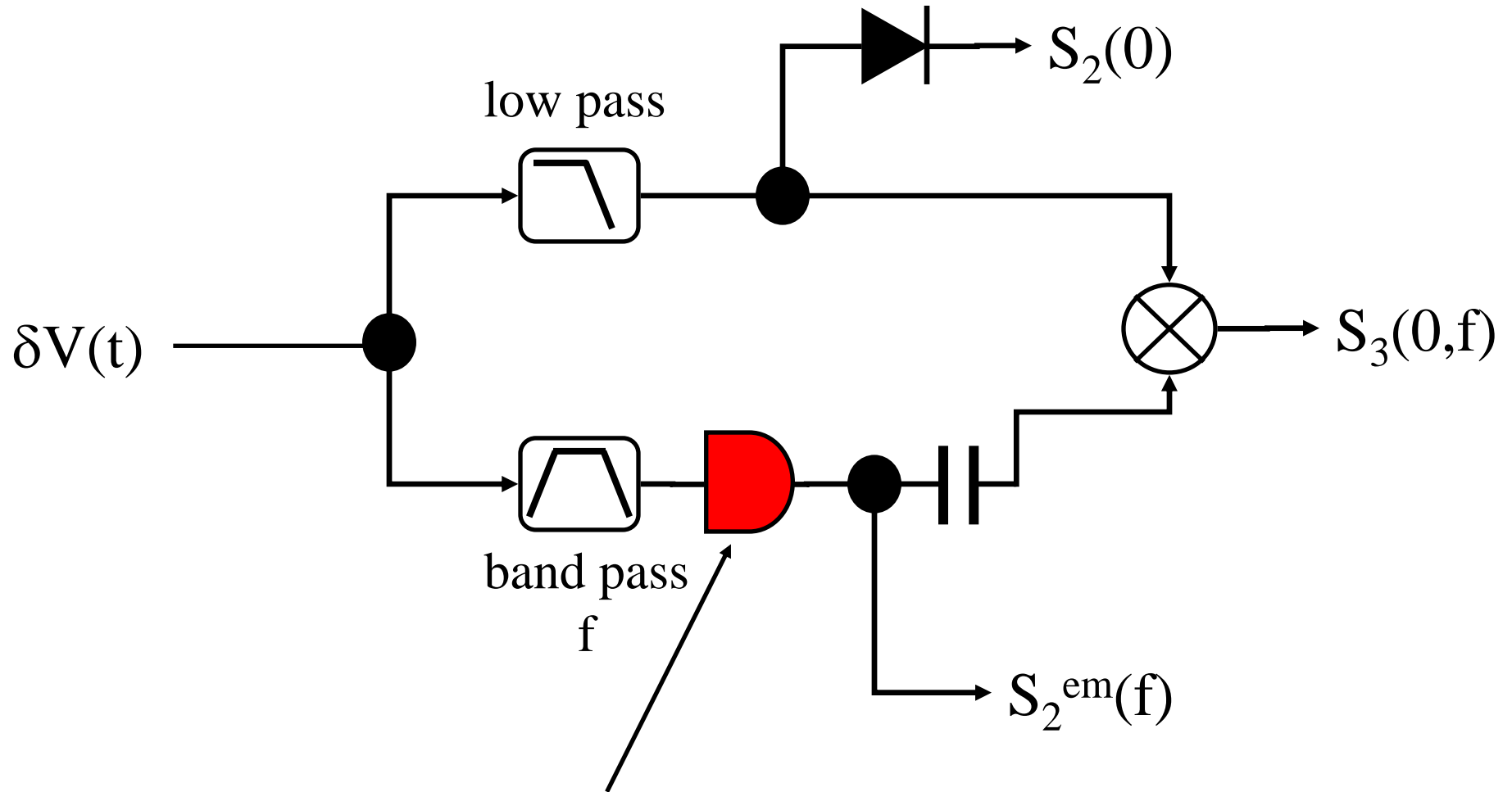
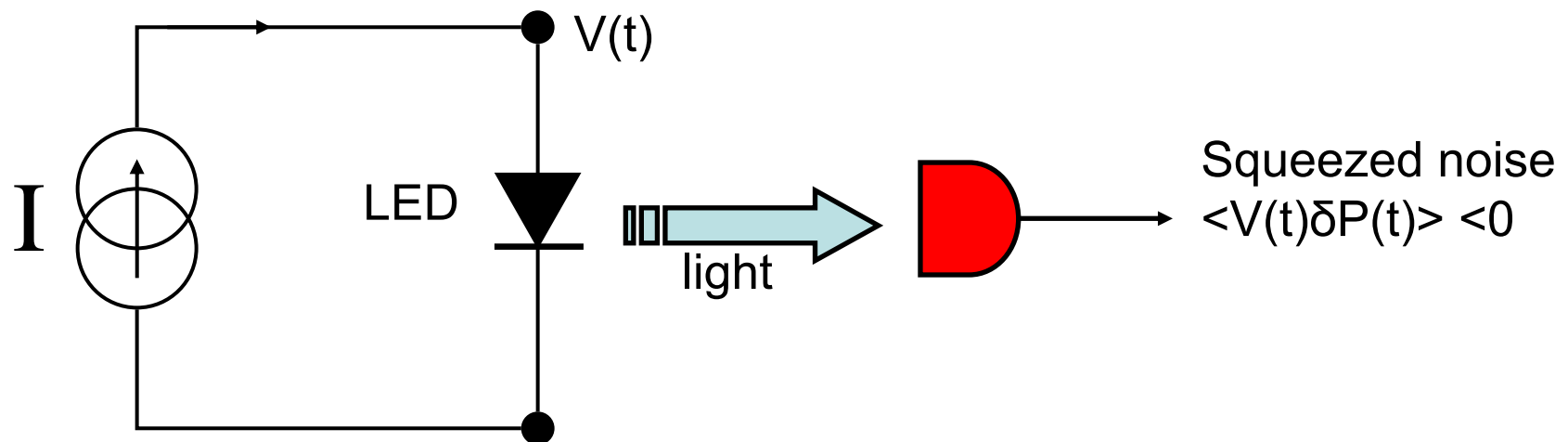
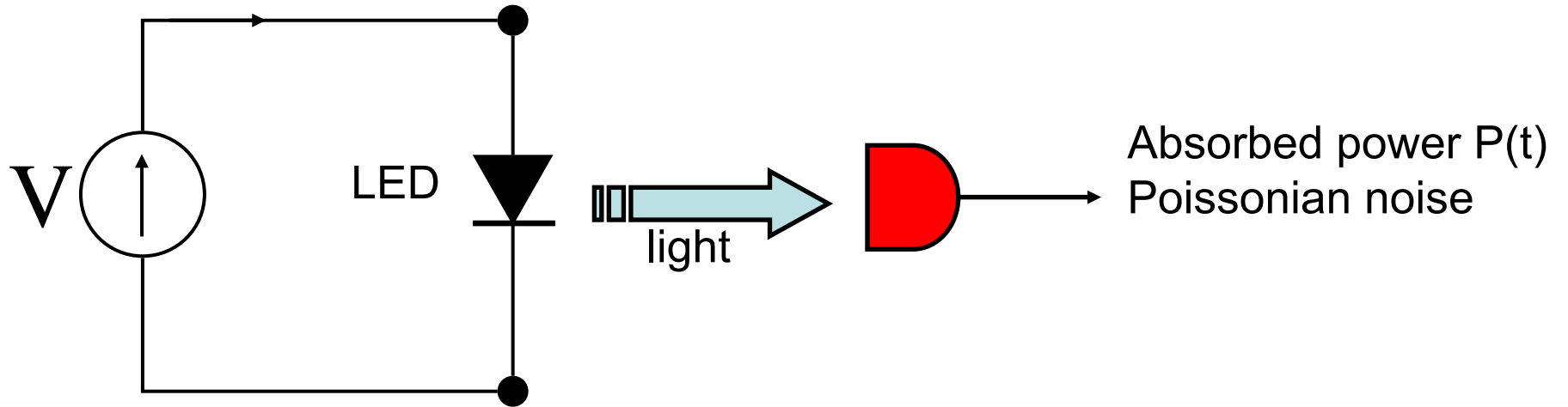


Photo-multiplier: absorbs photons

Similar Q optics experiments...



Conclusion

NOISE SUSCEPTIBILITY

Probe for inelastic mechanisms:

- * electron-phonon time vs T in normal metals at very low T
- * electron-electron time
- * in more subtle systems (carbon nanotubes, quasi-crystals, etc.)

Link with dynamical Coulomb blockade:

- * Van der Waals interaction between noise sources ?
- * Coulomb blockade at room T ?

THIRD CUMULANT OF NOISE

Test for deep understanding of transport mechanisms ?

Measurement theory: how to order three current operators in quantum mechanics ?

Which system exhibits gaussian noise ?

- * Not a radiator, not a light bulb run through by a current (more current, more heat or light)...
- * Not the cosmologic background (?)
- * Not a chemical reaction (cf tunnel junction)
- * Not a finite system, even at equilibrium (there is S_4 ...)

ONLY a perfectly transmitting channel...